Early Results from GLASS-JWST. X. Rest-frame UV-optical Properties of Galaxies at $7 < z < 9$

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Abstract

We present the first James Webb Space Telescope/NIRCam-led determination of $7 < z < 9$ galaxy properties based on broadband imaging from 0.8 to 5 μm as part of the GLASS-JWST Early Release Science program. This is the deepest data set acquired at these wavelengths to date, with an angular resolution $\lesssim 0''\!.14$. We robustly identify 13 galaxies with signal-to-noise ratio $\gtrsim 8$ in F444W from 8 arcmin^2 of data at $m_{AB} \leqslant 28$ from a combination of dropout and photometric redshift selection. From simulated data modeling, we estimate the dropout sample purity to be $\gtrsim 90\%$. We find that the number density of these F444W-selected sources is broadly consistent with expectations from the UV luminosity function determined from Hubble Space Telescope data. We characterize galaxy physical properties using a Bayesian spectral energy distribution fitting method, finding a median stellar mass of $10^{8.5} M_\odot$ and age 140 Myr, indicating they started ionizing their surroundings at redshift $z > 9.5$. Their star formation main sequence is consistent with predictions from simulations. Lastly, we introduce an analytical framework to constrain main-sequence evolution at $z > 7$ based on galaxy ages and basic assumptions, through which we find results consistent with expectations from cosmological simulations. While this work only gives a glimpse of the properties of typical galaxies that are thought to drive the reionization of the universe, it clearly shows the potential of JWST to unveil unprecedented details of galaxy formation in the first billion years.

Unified Astronomy Thesaurus concepts: Galaxy evolution (594); High-redshift galaxies (734)

Supporting material: machine-readable tables

1. Introduction

Redshift $z \sim 8$–7 is likely the epoch when the universe undergoes a rapid phase transition. The intergalactic medium neutral fraction $\bar{x}_{HI}$ quickly declines from $\bar{x}_{HI} \sim 1$ to ~ 0 with decreasing redshift (e.g., Treu et al. 2013; McGreer et al. 2015; Greig et al. 2017; Bañados et al. 2018; Mason et al. 2018). For a full understanding of this complex reionization process, it is fundamental to characterize the physical properties of the sources, especially galaxies around and below the characteristic luminosity L_*, as those are thought to contribute the most to the photon budget due to their high abundance (e.g., Robertson et al. 2015). Key topics to study include, but are not limited to, their star formation rates, underlying stellar population properties such as

mass, ages, and metallicities, as well as the morphology and dust content.

Spectral energy distribution (SED) fitting is a crucial process to derive galaxy physical parameters (e.g., Maraston et al. 2010; Roberts-Borsani et al. 2021). So far, the combination of Hubble and IRAC/Spitzer has been the primary tool for the community to determine the SED properties of $z \gtrsim 7$ galaxies. However, IRAC photometry is limited in terms of both depth and spatial resolution, which can lead to systematic uncertainties from deblending procedures (Merlin et al. 2015; Tacchella et al. 2022). The deepest data set on Spitzer/IRAC have 5σ depths of ~ 26–27 mag in the 4.5 μm band (e.g., Labbé et al. 2015; Stefanon et al. 2022a), which limits the current measurements to bright massive galaxies (stellar masses $> 10^9 M_\odot$, e.g., Roberts-Borsani et al. 2022). For lower-mass galaxies, the measurements are mainly based on stacked SEDs (e.g., Stefanon et al. 2022b) or a handful of gravitationally lensed sources (e.g., Castellano et al. 2016; Atek et al. 2018; Bradač et al. 2019; Strait et al. 2021).

Moreover, given Spitzer's limited angular resolution, most of the existing $z \gtrsim 7$ galaxies are selected via detection in rest-frame UV bands either from Hubble Space Telescope (HST) images (e.g., Labbé et al. 2010; Roberts-Borsani et al. 2016) or from ground-based surveys (e.g., Whitler et al. 2022). Interestingly, ALMA has recently been playing an important role in high-z studies, yielding unprecedented observations of galaxies that are dark in HST bands (Franco et al. 2018; Shu et al. 2022). These results show that the selection of high-z galaxies based on rest-frame UV is partially biased and that redder wavelengths can potentially uncover an undiscovered population. The selection of $z > 7$ galaxies from the rest-frame optical band (i.e., $\sim 4\,\mu m$ observer frame) has not been possible so far, but the situation has now changed with the availability of James Webb Space Telescope (JWST hereafter) NIR-Cam data.

In this Letter, we present F444W-selected F090W-dropout candidates ($7 \lesssim z \lesssim 9$) in the NIRCam parallel pointing from the JWST GLASS Early Release Science Program (Treu et al. 2022a), the deepest image at these wavelengths obtained so far from the new observatory. This is effectively approximately a rest-frame V-band selection. We report the physical properties of the identified candidates, such as luminosities, stellar masses, star formation rates (SFRs), and ages. Further, we use the galaxy properties to constrain the star formation main sequence at this redshift and above. With the sensitivity of JWST and depth of the ERS program, we investigate galaxies with $10^8 \lesssim M_*/M_\odot \lesssim 10^9$ (F444W $\lesssim 28$ mag), approaching a mass range that likely includes many of the sources that are responsible for reionization (e.g., Alvarez et al. 2012).

We quote magnitudes in the AB system (Oke & Gunn 1983) and adopt the standard cosmology with $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$, and $H_0 = 70\,km\,s^{-1}\,Mpc^{-1}$.

2. Data and Sample Selection

We use the photometric catalog of the first NIRCam parallel observations of the GLASS-JWST-ERS program obtained on 2022 June 28–29 (Paper II; Merlin et al. 2022). The observation was acquired in parallel to the primary NIRISS observation of cluster Abell 2744. The field is sufficiently far to avoid significant gravitational-lensing magnification (approximately one virial radius away from the cluster core). Although magnification at the $\sim 30\%$ level is still expected based on weak-lensing studies (Medezinski et al. 2016), we assume that all galaxies in this work have magnification $\mu = 1$. The data set includes seven wide filters: F090W, F115W, F150W, F200W, F277W, F356W, and F444W. Their 5σ depths in a $0.''30$ diameter aperture are in the range of 29–29.5 AB mag. We refer readers to Paper II for the detailed description of the image reduction, sensitivity, and photometric catalog construction. The final images are PSF-matched to the F444W band (coarsest resolution with FWHM $\approx 0.''14$). Objects are detected with SourceExtractor (Bertin & Arnouts 1996) using F444W as the detection image. That is, for $z \sim 7$–8, the galaxies are detected based on the rest-frame V band. In this work, the colors refer to magnitudes measured in circular apertures of size $0.''30$ in diameter (~ 2 FWHM). The total F444W fluxes are the Kron fluxes measured with A-PHOT (Merlin et al. 2019), taken directly from the catalog in Paper II. For other bands, the total fluxes are $0.''30$ diameter aperture fluxes scaled with the ratios between the total fluxes and the $0.''30$ diameter aperture fluxes in the F444W band. Our aperture size choice is

motivated so that it captures most of the fluxes from the galaxies at these redshifts while minimizing the risk of contamination from other nearby sources. As measured in Yang et al. (2022, Paper V), our final galaxy candidates have a median effective radius of 0.54 kpc in the F444W band, which corresponds to $0.''10$ at $z = 8$). We note if we were to use $0.''45$ diameter aperture fluxes to scale total flux calculations, our results would remain substantially unchanged within the limits of photometric scatter uncertainties.

2.1. Color and Signal-to-noise Ratio Selection Criteria

We optimize the selection criteria by following the framework described in Hainline et al. (2020) based on the JAdes extraGalactic Ultradeep Artificial Realizations (JAGUAR) mock catalog (Williams et al. 2018). We aim to obtain a high-purity F090W dropout sample, which led to the following criteria:

$$
\begin{aligned}
&\mathrm{F090W} - \mathrm{F115W} > 0.75 \\
&\mathrm{F150W} - \mathrm{F200W} < 0.4 \\
&\mathrm{F090W} - \mathrm{F115W} > 1.5 \times (\mathrm{F150W} - \mathrm{F200W}) + 1 \\
&\quad \mathrm{S/N\,(F090W)} < 2.0
\end{aligned} \tag{1}
$$

We also require that the signal-to-noise ratio (S/N) be >8 in the F444W band and >2 in the F115W, F150W, and F200W bands. We note that if sources satisfy the selection criteria for F115W dropouts ($z > 9$) as presented in Castellano et al. (2022, Paper III), they are included in that sample instead.

A four-band color selection is recommended by Hainline et al. (2020) as it yields higher completeness and accuracy than a three-band color selection. Generally, the nondetection criteria are only applied to bands below the dropout band. However, in our simulated data analysis, we find that the requirement of S/N (F090W) < 2 rejects $\sim 75\%$ of $z \sim 2$ galaxies that would otherwise have leaked in as interlopers. In addition, the S/N (F090W) criterion helps reject $\sim 70\%$ of $6 < z < 7$ sources that would have otherwise been included in the sample. Including the F090W S/N criterion has a minimal effect on the candidates with true redshift $z > 7.25$ and merely shifts the average redshift of the F090W dropouts from 7.2 to 7.5.

Based on these JAGUAR tests, the color and S/N selections yield a sample with purity $\gtrsim 95\%$ for objects brighter than F444W $\leqslant 27.5$ mag. Here we define purity as a ratio of the number of candidates with true redshift $z > 7$ to the number of all selected candidates. The purity drops quickly to $\sim 50\%$ at F444W $= 28.5$ mag, where the contamination by $z \sim 2$ galaxies is about 30%. The other 20% of the objects that passed the selection criteria are $5 < z < 7$ galaxies. Finally, we test the selection criteria on the mock NIRCam images of the observed field described in Paper II. We find a similar conclusion in terms of both purity and completeness.

2.2. Photometric Redshift Refinement

To further refine the sample selection, we follow Kauffmann et al. (2020), who found that adding a photometric redshift (photo-z) selection to a color cut yields higher purity with minimal loss of completeness. Thus, we refine the color–S/N-selected sample with the photometric redshift code EAzY (Brammer et al. 2008) with the default V1.3 spectral template and a flat prior.

Based on the output photometric redshifts, we restrict the selection to candidates (1) with the most probable redshift at high z and (2) with an integrated probability under the peaks of
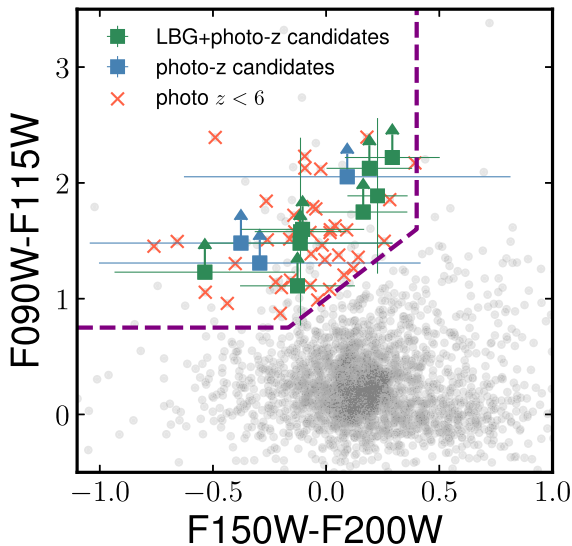
Figure 1. Observed color selection diagram for F090W dropouts. Green squares indicate the LBG+photo-z candidates in this paper (see Section 2.1). Orange crosses are objects that pass the color and S/N selection criteria but failed the photo-z criterion. The gray dots show all objects in the photometric catalog that pass the minimum S/N criteria (i.e., S/N(F444W) > 10 and S/N (F115W, F150W, F200W) > 2). The gray dots that are in the color selection box are those with S/N(F090W) > 2 and do not pass the S/N(F090W) criterion. Blue squares are photo-z candidates (see Section 2.3).

the distribution that is largest at high z, i.e., $z_{\rm peak} > 6$ and $z_p > 6$ in EAzY nomenclature. In the test on the JAGUAR mock catalog, this procedure reduces the interloper contamination from ∼30% to ≲10% at F444W = 28.5 mag. The completeness drops by an additional few percent for sources with F444W brighter than 27 mag, by ∼20% at 28 mag, and by ∼40% at 28.5 mag.

As for candidates from the actual JWST images, we have a total of 55 candidates from the color and S/N cut alone. We plot the color selection diagram of these galaxies in Figure 1. The photo-z refinement rejects 41 objects. Among those that fail the photo-z criterion, 25 are either artifacts, partially fall into gaps, or blobs with visible sizes larger than $1''$. For the rest, most have fluxes in the F090W band with S/N > 1 and are fainter than F444W > 27 mag, except for three sources that could well be at $z > 7$ based on colors and visual inspection but are still excluded from further analysis for consistency of the selection criteria. As a result, we have 14 candidates for further refinement. We call this type of candidates the LBG+photo-z sample. They are our primary sample.

2.3. Additional Photometric Candidates

In addition to the LBG+photo-z sample above, we perform a second independent selection to capture additional high-z candidates that fail to enter into the strict LBG selection window. This additional sample may include the sources within the color–color region copopulated by high-redshift sources and passive intermediate-redshift interlopers, and the sources that may be faint in the F115W–F200W bands. To do so, we carry out a photometric redshift analysis of all objects in the source catalog and select galaxies that have best-fit photometric redshifts in the range of $7 < z < 9$ assessed by two photometric redshift codes: EAzY and zphot (Fontana et al. 2000). We refer readers to Santini et al. (2022, Paper XI) for further details

of the zphot fitting procedure. Furthermore, we require that their $0''.1$ central regions must not fall in gaps in at least six out of the seven bands, including both F090W and F115W, which we deem critical as those bands encompass the Lyman break. Lastly, at least three of the bands from F115W to F356W must have S/N > 2, and F444W must have S/N ≳ 8. Based on mock-catalog analysis, having photometric redshift confirmed by two independent codes increases the purity of the sample. Based on simulated NIRCam images of the observed field, we estimate a sample purity of >70% for this additional photometric selection.

We find 14 candidates, 9 of which are already included in the LBG+photo-z candidates. Therefore, we have five remaining objects for further refinement. We identify these objects as the photo-z sample.

2.4. Contamination by Cool Dwarfs

Although the expected number of ultracool dwarfs is only a few per field (Ryan & Reid 2016), we take a precautionary step to remove possible contamination from late L and T dwarfs in our sample. One effective solution to mitigate this contamination, without resorting to the characterization of the light profile of the sources to determine whether they are spatially resolved or point sources, is to exclude objects with F356W − F444W < −0.3 (Hainline et al. 2020). With this criterion, we reject two additional objects from the LBG+photo-z sample and none from the photo-z sample.

2.5. Visual Inspection

Finally, two authors (N.L. and M.T.) visually reviewed all candidates. The objective of this process is strictly limited to eliminating defects and artifacts that could escape automatic flagging, such as spikes of foreground stars, scattered light, and objects that fall into gaps. Such a sanity check is especially important for a new camera and telescope. To be conservative, we only include objects that both reviewers independently approved. Here we remove three objects from the LBG+photo-z sample. One is an artifact known as "dragon's breath" in the JWST artifact nomenclature (Rigby et al. 2022). One partially falls in a gap in F090W, and one is a large faint blob next to bright objects. We remove one object from the photo-z sample that partially falls in the gap in the F090W image.

3. Results and Discussion

We obtain nine final LBG+photo-z candidates and four final photo-z candidates at 8σ detection in F444W. Their image stamps are shown in Figures A1 and A2. We discuss their photometric properties in the Appendix. The F444W magnitude range of the candidates is 26.3–28.0 with a median of 27.6 mag. We discuss their SED-inferred physical properties in the sections below. Other properties such as mass-to-light ratios, morphology, and sizes are discussed in detail in Paper XI, Treu et al. (2022b, Paper XII), and Paper V, respectively.

The total area where the bands required for candidate selections (F090W, F115W, F150W, F200W, and F444W) overlap is 7.7 arcmin^2. Based on the segmentation map, 87% are not covered by foreground objects. Our preliminary test on mock catalogs suggests that the completeness of the LBG +photo-z selection is ∼40%. Using the existing UV luminosity functions (Bouwens et al. 2015; Bowler et al. 2020), we expect ∼10 LBG+photo-z candidates at $7 < z < 9$ in the field, which

is in agreement with the number of candidates found here. Since previous UV luminosity functions were mainly based on rest-frame UV-selected galaxies, our results suggest that the population of LBG galaxies that are selected from the rest-frame V band are similar to those selected from the rest-frame UV band. Nonetheless, this comparison is only qualitative because we do not have detailed source recovery and completeness simulations available for this preliminary version of the image reduction.

3.1. SED Modeling

We use the Bayesian SED-fitting code `Bagpipes` (Carnall et al. 2018) to infer the physical properties of the galaxies in our sample. The code is based on the 2016 version of the Bruzual & Charlot (2003) stellar population models, the `CLOUDY` photo-ionization code (Ferland et al. 2017), the Inoue et al. (2014) IGM attenuation model, and the Kroupa & Boily (2002) IMF. We assume the Calzetti et al. (2000) dust attenuation law. The free parameters and their prior ranges that are not related to SFH are stellar mass ($6 < \log(M_*/M_\odot) < 13$), attenuation in the V band ($0 < A_V < 3$), ionization parameter ($-4 < \log U < -2$), and metallicity ($0 < Z/Z_\odot < 1$).

While parameters derived from SED fitting are relatively independent of model templates (Whitler et al. 2022), fitting codes, and assumed dust attenuation law (Topping et al. 2022), they are dependent on the assumed SFH (Carnall et al. 2019; Leja et al. 2019; Topping et al. 2022; Whitler et al. 2022). The dependence is stronger for the age and weaker for the stellar mass and SFR (Santini et al. 2015; Tacchella et al. 2022). For this reason, we adopt a log-normal SFH, which allows a wide range of SFH shapes, ranging from rising to declining over time, despite its parametric nature. Moreover, this functional form matches well the overall shapes of the majority of SFHs in cosmological hydrodynamic simulations (e.g., Gladders et al. 2013; Diemer et al. 2017). The free parameters under this assumption are the age of the universe at the peak of star formation t_{peak} and the FWHM of the SFH duration. We set both to vary between 1 Myr and 13 Gyr.

For each object, we run `Bagpipes` twice. First, with a redshift fixed to the z_{peak} redshift from `EAZY`. The best-fit values of the inferred physical properties quoted in this paper are from this step. The redshift is fixed to the `EAZY` best-fit values because (1) `EAZY` is a common photometric redshift SED-fitting tool that has been used and validated in multiple deep surveys (e.g., Skelton et al. 2014; Straatman et al. 2016) and (2) for consistency with other papers in the series that use `EAZY` output as the reference photometric redshifts, e.g., Glazebrook et al. (2022, Paper XV) and Nanayakkara et al. (2022, Paper XVI). For the second run of `Bagpipes`, we allow the redshift to vary within 3σ of z_{peak}. The uncertainties reported in this paper account for redshift uncertainties from this step. The best-fit SEDs at z_{peak} are shown in Figure 2. We list the inferred physical properties in Table 1.

The stellar mass of our candidates ranges from $10^{8.0}$ to $10^{9.3} M_\odot$ with a median of $10^{8.5} M_\odot$. The mass-weighted age ranges from 35 to 220 Myr with a median of 140 Myr. The ages correspond to a formation redshift (the redshift at which half of the current stars in each galaxy have formed) of $z = 10.6$ to $z = 7.8$, with a median of $z = 9.5$. This suggests that the majority of these galaxies started forming stars before redshift 9.5. For the more massive galaxies with ages close to 200 Myr, we infer substantial star formation happening at $z \gtrsim 11$, which

is consistent with the $z > 11$ objects found in the same data set (Paper III; Naidu et al. 2022). The inferred dust extinction (A_V) has an average of $0.2^{+0.2}_{-0.1}$ mag, which is consistent with $A_V \sim 0.5$ mag of the $z = 8.38$ A2744_YD4 lensed galaxy (Laporte et al. 2017), as modeled by Behrens et al. (2018). This average extinction is a few times lower than typical values at $z \sim 2$–3 (e.g., Theios et al. 2019).

Our best-fit SFHs all show SFRs that increase with time, in approximate agreement with what is found in cosmological simulations at similar mass and redshift without burstiness features (e.g., Ma et al. 2018). As a cautionary note, based on Whitler et al. (2022), about $\lesssim 10\%$ of the $z \sim 7$ galaxies undergo an intense burst of star formation. For these galaxies, the assumed smooth SFH would be skewed toward the recent burst, likely resulting in an underestimation of mass and age.

3.2. SFR of $7 < z < 9$ Galaxies

We plot the star formation main sequence in the left panel of Figure 3. Our measurements are shown in colors. Gray data points show the measurements in literature based on SED modeling with nebular emission lines. We fit our measurements with a linear relation in log scale using a least-squares method and Monte Carlo sampling (repeated 1000 times):

$$\log_{10}\left(\frac{\text{SFR}}{M_\odot \text{ yr}^{-1}}\right) = \alpha \log_{10}\left(\frac{M_*}{10^8 M_\odot}\right) + \beta, \qquad (2)$$

where α and β are free parameters. The best-fit slope and normalization at $10^8 M_\odot$ are $0.76^{+0.15}_{-0.14}$ and $-0.08^{+0.10}_{-0.14}$, with an average scatter of 0.1 dex.

We follow the method in Santini et al. (2017) to correct the fitted linear equation for Eddington bias. In short, we replicate the observed mass 20 times to increase the statistics and consider these as "true" masses. We then create a grid of true α's and β's. For each "true" α and β, we calculate each galaxy's "true" main-sequence SFR with the observed scatter. We then add Gaussian noise based on each object's measurement uncertainties to obtain the "observed" mass and SFR. We calculate the "observed" best-fit linear fit for that "true" α and β. The Eddington-bias corrected parameters are the range of these true α and β that match the observed slope and normalization and their uncertainty ranges. We find that the Eddington-bias corrected parameters are $\alpha = 0.95^{+0.17}_{-0.23}$ and $\beta = -0.15^{+0.13}_{-0.12}$.

The obtained main-sequence slope α is consistent with lower redshifts (Santini et al. 2017). Both hydrodynamical simulations and semianalytical models predict α that is close to unity, e.g., $\alpha = 1.03$ in Ma et al. (2018) and $\alpha = 0.95$ in Yung et al. (2019). Both are consistent with our measurement.

3.3. SFR Evolution Insight from Galaxy Ages

Age as a function of stellar mass provides additional information to constrain the relation between star formation and mass at even higher redshifts, where it might be difficult to obtain large-enough samples from direct observations, even with JWST. Here we present a method to constrain the time-dependent component in the main-sequence evolution. We assume that SFR ($\dot{M}$) is a power-law function of both stellar mass and redshift:

$$\dot{M} \propto M^\alpha (1 + z)^\gamma \propto M^\alpha (t)^{\gamma'}. \qquad (3)$$

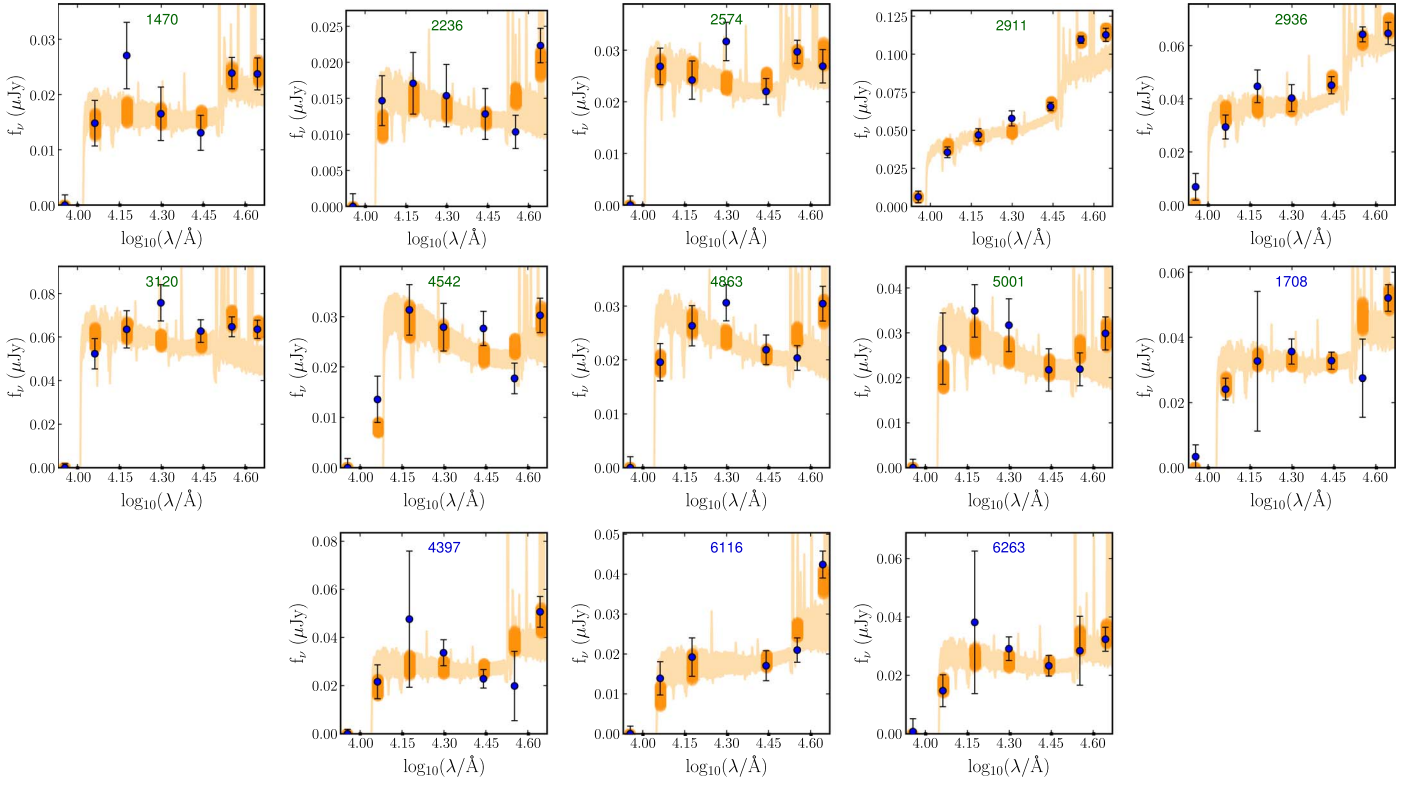M and t are the stellar mass and time since the big bang. In the matter-dominated universe (which is relevant to the

Figure 2. Best-fit SEDs of $7 \lesssim z \lesssim 9$ candidates, derived with the most probable redshift (z_{peak}) determined by EAzY. Observed total fluxes are shown in blue. The first nine objects with IDs in green are the LBG+photo-z candidates. The last four objects with IDs in blue are the photo-z candidates.

Table 1
The Physical Properties Inferred from the F090W Dropout Galaxies

ID	R.A.	Decl.	z(EAzY)	$\log(M_*/M_\odot)$	SFR ($M_\odot$/yr)	Age (Myr)	M_{UV}[a]	A_V
			LBG+photo-z candidates					
1470	3.5137326	−30.3628201	$7.6^{+0.5}_{-0.5}$	$8.4^{+0.2}_{-0.1}$	$1.4^{+0.4}_{-0.2}$	160^{+76}_{-55}	$-18.7^{+0.2}_{-0.3}$	$0.2^{+0.2}_{-0.1}$
2236	3.4898570	−30.3544532	$8.0^{+0.6}_{-0.7}$	$8.0^{+0.3}_{-0.7}$	$0.9^{+0.3}_{-0.6}$	72^{+88}_{-67}	$-18.7^{+0.3}_{-0.1}$	$0.1^{+0.2}_{-0.1}$
2574	3.4954699	−30.3507722	$7.4^{+0.3}_{-0.3}$	$8.4^{+0.1}_{-0.3}$	$1.6^{+0.3}_{-0.3}$	140^{+48}_{-95}	$-19.2^{+0.1}_{-0.2}$	$0.1^{+0.1}_{-0.1}$
2911	3.5118088	−30.3468413	$6.9^{+0.2}_{-0.2}$	$9.3^{+0.2}_{-0.1}$	$9.1^{+3.6}_{-1.9}$	180^{+81}_{-69}	$-19.5^{+0.1}_{-0.2}$	$0.7^{+0.1}_{-0.1}$
2936	3.5119322	−30.3466618	$7.2^{+0.6}_{-5.0}$	$9.1^{+0.1}_{-0.2}$	$4.2^{+1.5}_{-0.9}$	230^{+74}_{-120}	$-19.4^{+0.1}_{-0.5}$	$0.4^{+0.2}_{-0.2}$
3120	3.5202581	−30.3439199	$7.4^{+0.3}_{-0.3}$	$8.5^{+0.3}_{-0.2}$	$3.2^{+1.3}_{-1.3}$	35^{+71}_{-22}	$-20.1^{+0.1}_{-0.3}$	$0.2^{+0.1}_{-0.2}$
4542	3.4879853	−30.3254241	$9.0^{+0.7}_{-0.6}$	$8.5^{+0.2}_{-0.3}$	$2.0^{+0.4}_{-0.5}$	110^{+42}_{-63}	$-19.5^{+0.2}_{-0.1}$	$0.1^{+0.1}_{-0.1}$
4863	3.4866961	−30.3272162	$8.1^{+0.4}_{-0.5}$	$8.1^{+0.3}_{-0.3}$	$1.4^{+0.4}_{-0.7}$	52^{+64}_{-35}	$-19.3^{+0.1}_{-0.1}$	$0.1^{+0.1}_{-0.0}$
5001	3.4997064	−30.3177257	$8.1^{+0.7}_{-0.8}$	$8.3^{+0.2}_{-0.4}$	$1.6^{+0.4}_{-0.7}$	74^{+67}_{-52}	$-19.4^{+0.2}_{-0.1}$	$0.1^{+0.1}_{-0.0}$
			photo-z candidates					
1708	3.4905553	−30.3603869	$7.8^{+0.4}_{-0.6}$	$8.8^{+0.2}_{-0.2}$	$3.3^{+0.8}_{-0.6}$	150^{+68}_{-79}	$-19.5^{+0.1}_{-0.3}$	$0.3^{+0.1}_{-0.2}$
4397	3.4746591	−30.3226226	$8.1^{+0.6}_{-0.6}$	$8.8^{+0.2}_{-0.2}$	$2.6^{+1.3}_{-0.5}$	160^{+74}_{-82}	$-19.4^{+0.3}_{-0.2}$	$0.2^{+0.3}_{-0.1}$
6116	3.5045910	−30.3079217	$8.2^{+0.6}_{-0.6}$	$8.7^{+0.2}_{-0.7}$	$2.6^{+1.4}_{-1.4}$	140^{+96}_{-120}	$-18.8^{+0.4}_{-0.2}$	$0.5^{+0.3}_{-0.2}$
6263	3.4696869	−30.3090295	$8.2^{+0.5}_{-5.7}$	$8.5^{+0.2}_{-0.4}$	$2.1^{+0.7}_{-1.9}$	130^{+630}_{-74}	$-19.3^{+3.6}_{-0.2}$	$0.1^{+0.3}_{-0.1}$

Notes. We model the observed photometry with BAGPIPES assuming a log-normal SFH. The values are the median of the posterior distribution when redshift is fixed to z_{peak} from EAZY. The uncertainties include the marginalized 1σ interval when the redshift is free to vary (see text). The age displayed is the mass-weighted age. All galaxies are assumed to have magnification $\mu = 1$.

[a] M_{UV} magnitudes are measured by sampling best-fit SEDs from the posterior distributions of the best-fit parameters. They are magnitudes corresponding to the average flux in the 100Å range centering at the 1500 Å wavelength.

(This table is available in its entirety in machine-readable form.)

calculation here), $\gamma' = -\frac{2}{3}\gamma$. Since the mass-weighted age can be approximated as

$$\mathrm{age}(t_0) \approx \frac{\int_0^{t_0} \dot{M}(t_0 - t)dt}{\int_0^{t_0} \dot{M}dt}, \quad (4)$$

where t_0 is the age of the universe at the observed redshift, we can use Equation (3) to show that

$$\mathrm{age}(M) \propto M^{\frac{1-\alpha}{1+\gamma'}}. \quad (5)$$

Here we predict a linear relation between the log mass-weighted age and log stellar mass. With this formalism, we can
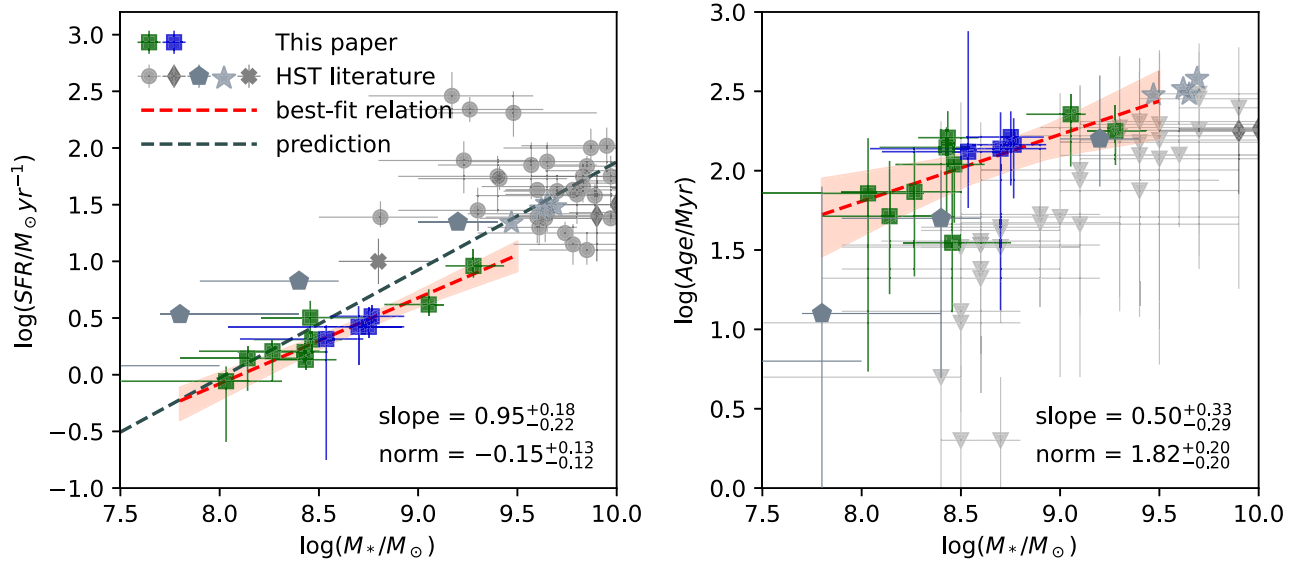
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Figure 3. Left: star formation main sequence (log SFR vs. log stellar mass). Right: log mass-weighted age vs. log stellar mass. Green and blue squares show our LBG +photo-z and photo-z candidates in this paper, respectively. The best-fit observed linear relation is plotted in red in each panel, while the Eddington-bias corrected slope and normalization at $10^8 M_\odot$ are listed at the lower right corners. The gray data points show literature values. Circles—rest-frame UV-selected $z \sim 7$ galaxies with ALMA FIR data (assumed nonparametric SFHs) from Topping et al. (2022). Diamond—$z \sim 8.5$ galaxies in CANDELS fields with Spitzer/IRAC data (assumed nonparametric SFHs) from Tacchella et al. (2022). Pentagons—stacked $z \sim 8$ galaxies in the HST legacy fields with deep Spitzer/IRAC data (assumed constant SFHs) from Stefanon et al. (2022b). Stars and filled cross—$z \sim 8$ galaxies with IRAC detections from SuperBoRG (assumed delayed-tau and nonparametric SFHs, respectively) (Morishita et al. 2020; Roberts-Borsani et al. 2022). Triangles—$z \sim 6.8$ galaxies in the COSMOS field with Spitzer/IRAC data (assumed constant SFHs) from Whitler et al. (2022). The gray dashed line indicates the theoretical prediction at $z = 8$ from Yung et al. (2019).

constrain the time-dependent term of the star formation rate γ' from its slope and the α term obtained in the above section. We note that this approach is approximate under the assumption that the galaxies have not started their quenching process and have been, on average, residing on the main sequence, i.e., no significant burst that produced a substantial fraction of stellar mass.

We plot log age versus log stellar mass in the right panel of Figure 3. We obtained a slope of $0.42^{+0.25}_{-0.25}$, a normalization of $1.81^{+0.18}_{-0.21}$, and a scatter of 0.2 dex. Using the same method to correct for the Eddington bias, we infer an intrinsic slope and normalization of $0.50^{+0.33}_{-0.29}$ and $1.82^{+0.20}_{-0.20}$. This indicates that the redshift-dependent term for the SFR evolution function is $\gamma = 1.4^{+0.7}_{-0.8}$. Our result is roughly consistent within 1σ with the predicted value of $\gamma = 2.01$ based on the FIRE-II simulation (Ma et al. 2018).

We note that ages measured on star-forming galaxies should always be treated as estimates only. Age is sensitive to multiple parameters such as the assumed SFH, metallicities, and binary fractions (e.g., Steidel et al. 2016; Leethochawalit et al. 2019). This is particularly true in galaxies with high specific star formation rate (Whitler et al. 2022). Such galaxies may include those with red F444W − F356W color, e.g., ID 4863 and ID 6116 in our sample, which are potentially due to high-equivalent-width nebular emission lines (Roberts-Borsani et al. 2016), indicating highly active star-forming activities.

4. Conclusion

We present 13 sources at $7 < z < 9$, photometrically selected from the NIRCam observation in the GLASS-JWST-ERS program. The galaxies have been identified based on rest-frame V-band detection, the first time that this has been possible for a sample of $L \lesssim L_*$ sources. Based on simulated data modeling, we expect our sample to have $\lesssim 10\%$ contamination. We

measured their stellar mass, SFRs, and mass-weighted ages using the `Bagpipes` SED-fitting code. The galaxies have stellar masses ranging from $10^{8.0}$ to $10^{9.3} M_\odot$. The key results from this work are:

1. Overall, the number density of these sources is broadly consistent with expectations from the UV luminosity function at this redshift determined from HST data. However, a detailed comparison is beyond the scope of this work as it requires exhaustive completeness and source recovery simulations.

2. The galaxies span a wide range in mass-weighted ages, from 30 Myr in the least massive galaxy to 100–200 Myr in the more massive $10^9 M_\odot$ galaxies. This suggests a significant amount of star formation in typical objects in our sample at redshift $z \gtrsim 11$. A caveat to note is that these conclusions are based on an assumed log-normal functional form for the SFH.

3. We present a preliminary determination of the star formation main sequence at $z \sim 8$. The Eddington-bias-corrected slope is $0.95^{+0.17}_{-0.23}$, consistent within 1σ with predictions from Ma et al. (2018) and Yung et al. (2019).

4. We present a method to constrain the star formation main-sequence evolution using inferred ages. We find that the redshift-dependent power-law slope term is $1.4^{+0.7}_{-0.8}$.

We note that the galaxies in this work may be affected by systematic uncertainties due to the preliminary NIRCam calibrations available at the time of writing. The fluxes were obtained at face values from the catalog in Paper II, where the uncertainties in the zero-points ($\sim$0.1 mag) are not included in the photometry calculations. Also, the field is affected by some modest amount of gravitational lensing, which would impact both mass and SFR measurements, but not age. The effect should be relatively minor for the star formation main sequence

since our inference is presented in a log scale. We will consider the lensing magnification in future work after the full data acquisition of the GLASS-ERS program.

Overall, this paper provides a first robust look at the properties of typical galaxies in the epoch of reionization from 19 hr of observations. It highlights the transformational capabilities of JWST infrared imaging to characterize the first light sources in the universe once both deeper and wider data sets become available.

Data Availability

We provide tables in this paper in machine-readable format. Further images, plots, and the routines used in this work are available upon reasonable request.

Appendix
Observed Photometric Properties of the F090W Dropouts

Figures A1 and A2 show the image stamps of our $7 < z < 9$ candidates in this work (the LBG+photo-z candidates and the photo-z candidates respectively). We list their photometric properties in Table A1. Their F444W magnitudes range from 26.3 to 28.0 with an average of $27.4^{+0.4}_{-0.5}$, while the F150W magnitudes range from 27.2 to 28.0 with an average of 27.6 ± 0.4 mag. At $7 < z < 8.5$, the F277W band is generally free of strong emission lines and can be used to probe the continuum below the 4000 Å break. However, Hβ + [O III] lines fall in the observed F444W band while the 4000Å break generally falls in the 356W band. There is no emission line-free band above the 4000Å break. The average F277W − F444W of our sample is $0.5^{+0.2}_{-0.3}$ mag, suggesting either the presence of the 4000Å break (and thus the presence of intermediate-age >0.3 Gyr stars) or strong Hβ + [O III] emission lines.

Table A1
Summary of Photometry of the F090W Dropout Galaxies

ID	F090W	F115W	F150W	F200W	F277W	F356W	F444W	m_{F444W}	Color1[a]	Color2[b]
					LBG+photo-z candidates					
1470	0.0 ± 1.8	14.8 ± 4.1	27.0 ± 6.0	16.5 ± 4.8	13.1 ± 3.2	23.8 ± 2.8	23.7 ± 2.9	28.0	−0.5	>1.2
2236	0.0 ± 1.8	14.6 ± 3.5	17.1 ± 4.3	15.3 ± 4.3	12.8 ± 3.5	10.3 ± 2.3	22.3 ± 2.4	28.0	−0.1	>1.5
2574	0.0 ± 1.8	26.8 ± 3.5	24.2 ± 3.7	31.6 ± 3.7	22.0 ± 2.5	29.6 ± 2.2	26.8 ± 3.2	27.8	0.3	>2.2
2911	6.2 ± 3.8	35.5 ± 3.5	46.9 ± 4.1	57.8 ± 4.9	65.6 ± 2.8	109.4 ± 2.4	112.6 ± 4.3	26.3	0.2	1.9
2936	6.9 ± 5.0	29.4 ± 4.5	44.7 ± 6.1	40.2 ± 5.0	45.0 ± 3.2	64.2 ± 2.8	64.6 ± 4.1	26.9	−0.1	1.6
3120	0.0 ± 1.9	52.1 ± 7.0	63.4 ± 8.6	75.6 ± 8.3	62.6 ± 5.2	64.6 ± 4.6	63.5 ± 4.2	26.9	0.2	>2.1
4542	0.0 ± 1.8	13.6 ± 4.6	31.3 ± 5.0	27.9 ± 4.7	27.6 ± 3.4	17.7 ± 3.0	30.2 ± 3.4	27.7	−0.1	>1.1
4863	0.0 ± 2.0	19.5 ± 3.5	26.3 ± 3.7	30.6 ± 3.4	21.8 ± 2.7	20.3 ± 2.3	30.4 ± 3.2	27.7	0.2	>1.7
5001	0.0 ± 1.9	26.5 ± 7.9	34.8 ± 5.8	31.6 ± 5.9	21.7 ± 4.7	21.9 ± 3.6	29.8 ± 3.7	27.7	−0.1	>1.6
					photo-z candidates					
1708	3.4 ± 3.6	24.1 ± 3.3	32.6 ± 21.4	35.6 ± 3.9	32.8 ± 2.6	27.4 ± 12.0	52.0 ± 4.1	27.1	0.1	>2.1
4397	0.0 ± 1.9	21.6 ± 7.0	47.6 ± 28.3	33.6 ± 5.4	22.8 ± 3.8	19.8 ± 14.4	50.6 ± 6.4	27.1	−0.4	>1.5
6116	0.0 ± 1.9	13.9 ± 4.2	19.2 ± 4.8	nan	17.0 ± 3.8	21.0 ± 3.0	42.4 ± 3.4	27.3	45.4	>1.1
6263	0.8 ± 4.4	14.7 ± 5.5	38.1 ± 24.4	29.1 ± 4.0	23.3 ± 3.5	28.4 ± 11.8	32.3 ± 4.1	27.6	−0.3	>1.3

Notes. The fluxes are total fluxes in nanojanskys. When the total flux is negative, we set the flux to zero in the SED-fitting procedure and set the uncertainty to the flux of the 1σ limiting magnitude. Colors are based on aperture magnitudes. For color calculation, the F090W aperture magnitudes are set to 1σ limiting magnitudes if the S/N is less than one and thus the F090W – F115W colors appear as lower limits. The IDs match those in Paper II.
[a] F150W − F200W UV continuum color.
[b] F090W − F115W Lyman drop color.

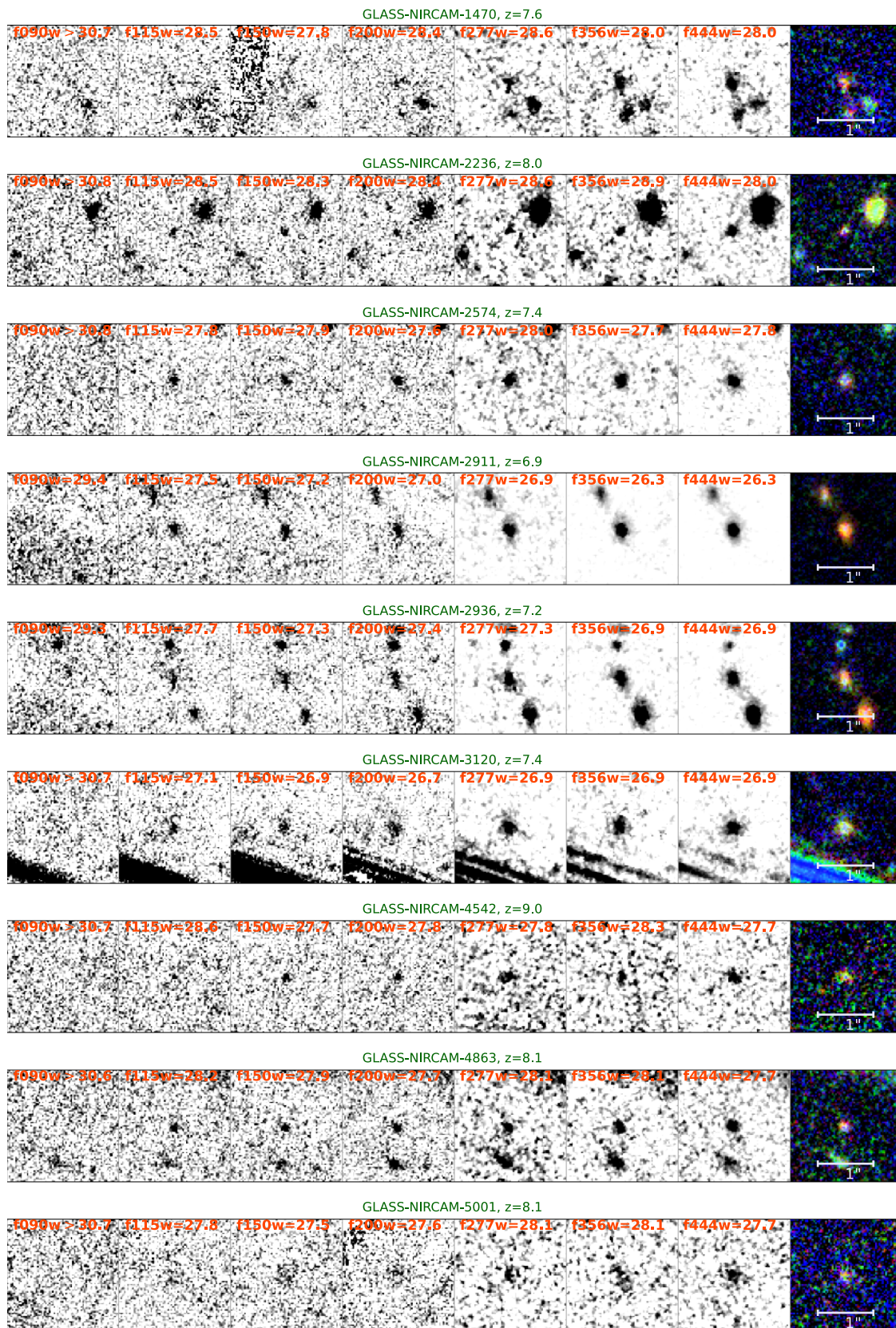(This table is available in its entirety in machine-readable form.)

Figure A1. Image stamps of $7 < z < 9$ LBG+photo-z candidates found in the GLASS-JWST-ERS NIRCam parallel observation at S/N(F444W) > 10. The columns from left to right are the images in F090W, F115W, F150W, F200W, F277W, F356W, F444W, and RGB composite based on the F444W, F277W, and F115W bands.
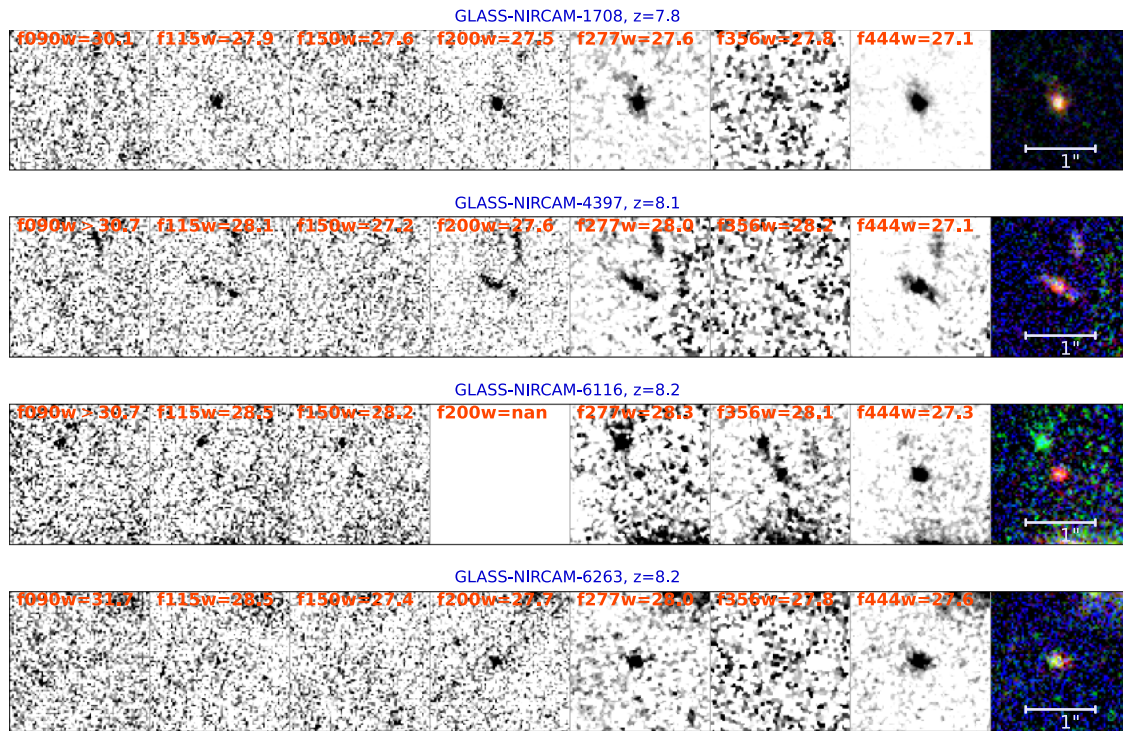
Figure A2. Same as Figure A1 but for the photo-z candidates.

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References

Alvarez, M. A., Finlator, K., & Trenti, M. 2012, ApJL, 759, L38
Astropy Collaboration, Price-Whelan, A. M., Sipőcz, B. M., et al. 2018, AJ, 156, 123
Atek, H., Richard, J., Kneib, J.-P., & Schaerer, D. 2018, MNRAS, 479, 5184
Bañados, E., Venemans, B. P., Mazzucchelli, C., et al. 2018, Natur, 553, 473
Behrens, C., Pallottini, A., Ferrara, A., Gallerani, S., & Vallini, L. 2018, MNRAS, 477, 552
Bertin, E., & Arnouts, S. 1996, A&AS, 117, 393
Bouwens, R. J., Illingworth, G. D., Oesch, P. A., et al. 2015, ApJ, 803, 34
Bowler, R. A. A., Jarvis, M. J., Dunlop, J. S., et al. 2020, MNRAS, 493, 2059
Bradač, M., Huang, K.-H., Fontana, A., et al. 2019, MNRAS, 489, 99
Brammer, G. B., van Dokkum, P. G., & Coppi, P. 2008, ApJ, 686, 1503
Bruzual, G., & Charlot, S. 2003, MNRAS, 344, 1000
Calzetti, D., Armus, L., Bohlin, R. C., et al. 2000, ApJ, 533, 682
Carnall, A. C., Leja, J., Johnson, B. D., et al. 2019, ApJ, 873, 44
Carnall, A. C., McLure, R. J., Dunlop, J. S., & Davé, R. 2018, MNRAS, 480, 4379
Castellano, M., Amorín, R., Merlin, E., et al. 2016, A&A, 590, A31
Castellano, M., Fontana, A., Treu, T., et al. 2022, ApJL, 938, L15
Diemer, B., Sparre, M., Abramson, L. E., & Torrey, P. 2017, ApJ, 839, 26
Ferland, G. J., Chatzikos, M., Guzmán, F., et al. 2017, RMxAA, 53, 385
Fontana, A., D'Odorico, S., Poli, F., et al. 2000, AJ, 120, 2206
Franco, M., Elbaz, D., Béthermin, M., et al. 2018, A&A, 620, A152
Gladders, M. D., Oemler, A., Dressler, A., et al. 2013, ApJ, 770, 64
Glazebrook, K., Nanayakkara, T., Jacobs, C., et al. 2022, arXiv:2208.03468
Greig, B., Mesinger, A., Haiman, Z., & Simcoe, R. A. 2017, MNRAS, 466, 4239
Hainline, K. N., Hviding, R. E., Rieke, M., et al. 2020, ApJ, 892, 125
Inoue, A. K., Shimizu, I., Iwata, I., & Tanaka, M. 2014, MNRAS, 442, 1805
Kauffmann, O. B., Le Fèvre, O., Ilbert, O., et al. 2020, A&A, 640, A67
Kroupa, P., & Boily, C. M. 2002, MNRAS, 336, 1188
Labbé, I., González, V., Bouwens, R. J., et al. 2010, ApJL, 716, L103
Labbé, I., Oesch, P. A., Illingworth, G. D., et al. 2015, ApJS, 221, 23
Laporte, N., Ellis, R. S., Boone, F., et al. 2017, ApJL, 837, L21
Leethochawalit, N., Kirby, E. N., Ellis, R. S., Moran, S. M., & Treu, T. 2019, ApJ, 885, 100
Leja, J., Carnall, A. C., Johnson, B. D., Conroy, C., & Speagle, J. S. 2019, ApJ, 876, 3
Ma, X., Hopkins, P. F., Garrison-Kimmel, S., et al. 2018, MNRAS, 478, 1694
Maraston, C., Pforr, J., Renzini, A., et al. 2010, MNRAS, 407, 830
Mason, C. A., Treu, T., Dijkstra, M., et al. 2018, ApJ, 856, 2
McGreer, I. D., Mesinger, A., & D'Odorico, V. 2015, MNRAS, 447, 499
Medezinski, E., Umetsu, K., Okabe, N., et al. 2016, ApJ, 817, 24
Merlin, E., Bonchi, A., Paris, D., et al. 2022, ApJL, 938, L14
Merlin, E., Fontana, A., Ferguson, H. C., et al. 2015, A&A, 582, A15
Merlin, E., Pilo, S., Fontana, A., et al. 2019, A&A, 622, A169
Morishita, T., Stiavelli, M., Trenti, M., et al. 2020, ApJ, 904, 50
Naidu, R. P., Oesch, P. A., van Dokkum, P., et al. 2022, ApJL, 940, L14
Nanayakkara, T., Glazebrook, K., Jacobs, C., et al. 2022, arXiv:2207.13860
Oke, J. B., & Gunn, J. E. 1983, ApJ, 266, 713

Rigby, J., Perrin, M., McElwain, M., et al. 2022, arXiv:2207.05632
Roberts-Borsani, G., Morishita, T., Treu, T., Leethochawalit, N., & Trenti, M. 2022, ApJ, 927, 236
Roberts-Borsani, G., Treu, T., Mason, C., et al. 2021, ApJ, 910, 86
Roberts-Borsani, G. W., Bouwens, R. J., Oesch, P. A., et al. 2016, ApJ, 823, 143
Robertson, B. E., Ellis, R. S., Furlanetto, S. R., & Dunlop, J. S. 2015, ApJL, 802, L19
Ryan, R. E. J., & Reid, I. N. 2016, AJ, 151, 92
Santini, P., Ferguson, H. C., Fontana, A., et al. 2015, ApJ, 801, 97
Santini, P., Fontana, A., Castellano, M., et al. 2017, ApJ, 847, 76
Santini, P., Fontana, A., Castellano, M., et al. 2022, arXiv:2207.11379
Shu, X., Yang, L., Liu, D., et al. 2022, ApJ, 926, 155
Skelton, R. E., Whitaker, K. E., Momcheva, I. G., et al. 2014, ApJS, 214, 24
Stefanon, M., Bouwens, R. J., Illingworth, G. D., et al. 2022a, ApJ, 935, 94
Stefanon, M., Bouwens, R. J., Labbé, I., et al. 2022b, ApJ, 927, 48
Steidel, C. C., Strom, A. L., Pettini, M., et al. 2016, ApJ, 826, 159
Straatman, C. M. S., Spitler, L. R., Quadri, R. F., et al. 2016, ApJ, 830, 51
Strait, V., Bradač, M., Coe, D., et al. 2021, ApJ, 910, 135
Tacchella, S., Finkelstein, S. L., Bagley, M., et al. 2022, ApJ, 927, 170
Theios, R. L., Steidel, C. C., Strom, A. L., et al. 2019, ApJ, 871, 128
Topping, M. W., Stark, D. P., Endsley, R., et al. 2022, MNRAS, 516, 975
Treu, T., Schmidt, K. B., Trenti, M., Bradley, L. D., & Stiavelli, M. 2013, ApJL, 775, L29
Treu, T., Roberts-Borsani, G., Bradac, M., et al. 2022a, ApJ, 935, 110
Treu, T., Calabro, A., Castellano, M., et al. 2022b, arXiv:2207.13527
Whitler, L., Stark, D. P., Endsley, R., et al. 2022, arXiv:2206.05315
Williams, C. C., Curtis-Lake, E., Hainline, K. N., et al. 2018, ApJS, 236, 33
Yang, L., Morishita, T., Leethochawalit, N., et al. 2022, ApJL, 938, L17
Yung, L. Y. A., Somerville, R. S., Popping, G., et al. 2019, MNRAS, 490, 2855